

# Einstein's equations with the Robertson-Walker Metric and perfect fluid energy-momentum tensor

=

$$\left(\frac{\dot{a}(t)}{a(t)}\right)^2 = \frac{8\pi G}{3}\rho(t) - \frac{k c^2}{a(t)^2}$$

$$\dot{\rho}(t) + 3\frac{\dot{a}(t)}{a(t)}\left(\rho(t) + \frac{p(t)}{c^2}\right) = 0$$

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$$\left(\frac{\dot{a}(t)}{a(t)}\right)^2 = \frac{8\pi G}{3}\rho(t) - \frac{kc^2}{a(t)^2}$$



$$H(t)^2 = \frac{8\pi G}{3}\rho(t) - \frac{kc^2}{a(t)^2}$$

$$H(t) \equiv \frac{\dot{a}(t)}{a(t)}$$

$$\dot{\rho}(t) + 3\frac{\dot{a}(t)}{a(t)}\left(\rho(t) + \frac{p(t)}{c^2}\right) = 0$$

# Critical density....?

$$H(t)^2 = \frac{8\pi G}{3}\rho(t) - \frac{kc^2}{a(t)^2}$$

$$\rho_{\text{crit}}(t) = \frac{3H(t)^2}{8\pi G}$$

$k=0$  gives critical density

$$\Omega_i(t) = \frac{\rho_i(t)}{\rho_{\text{crit}}(t)} = \rho_i(t) \frac{8\pi G}{3H(t)^2}$$

Now how do we go further...?

$$\dot{\rho}(t) + 3\frac{\dot{a}(t)}{a(t)}\left(\rho(t) + \frac{p(t)}{c^2}\right) = 0$$

➤ Need equation of state

$$\dot{\rho}_i(t) + 3\frac{\dot{a}(t)}{a(t)}\left(\rho_i(t) + \frac{p_i(t)}{c^2}\right) = 0$$

$$p/c^2 = \omega\rho$$

# Evolution of energy density as a function of a

$$\dot{\rho}_i(t) + 3\frac{\dot{a}(t)}{a(t)}\rho_i(t)(1+\omega) = 0$$

$$\frac{d\rho_i}{\rho_i} = -3(1+\omega)\frac{da}{a}$$

$$\rho_i(a) = \rho_{i,0} a^{-3(1+\omega)}$$

# Equation of state for:

↗ cold matter

$$\omega = 0$$

$$\rho_i(a) = \rho_{i,0} a^{-3(1+\omega)}$$

$$\rho_m \propto 1/a^3$$

↗ radiation

$$\omega = \frac{1}{3}$$

$$\rho_r \propto 1/a^4$$

↗ cosmological constant

$$\omega = -1$$

$$\rho_\Lambda \propto \text{const}$$

$$\rho c^2 = \frac{g}{(2\pi)^3} \int E(\mathbf{p}) f(\mathbf{p}) d^3|p|$$

$$p = \frac{g}{(2\pi)^3} \int \frac{|\mathbf{p}|^2}{3E} f(\mathbf{p}) d^3|p|$$

$$f(\mathbf{p}) = (\exp((E - \mu)/T) \pm 1)^{-1}$$

$$E^2 = m^2 c^4 + p^2 c^2$$

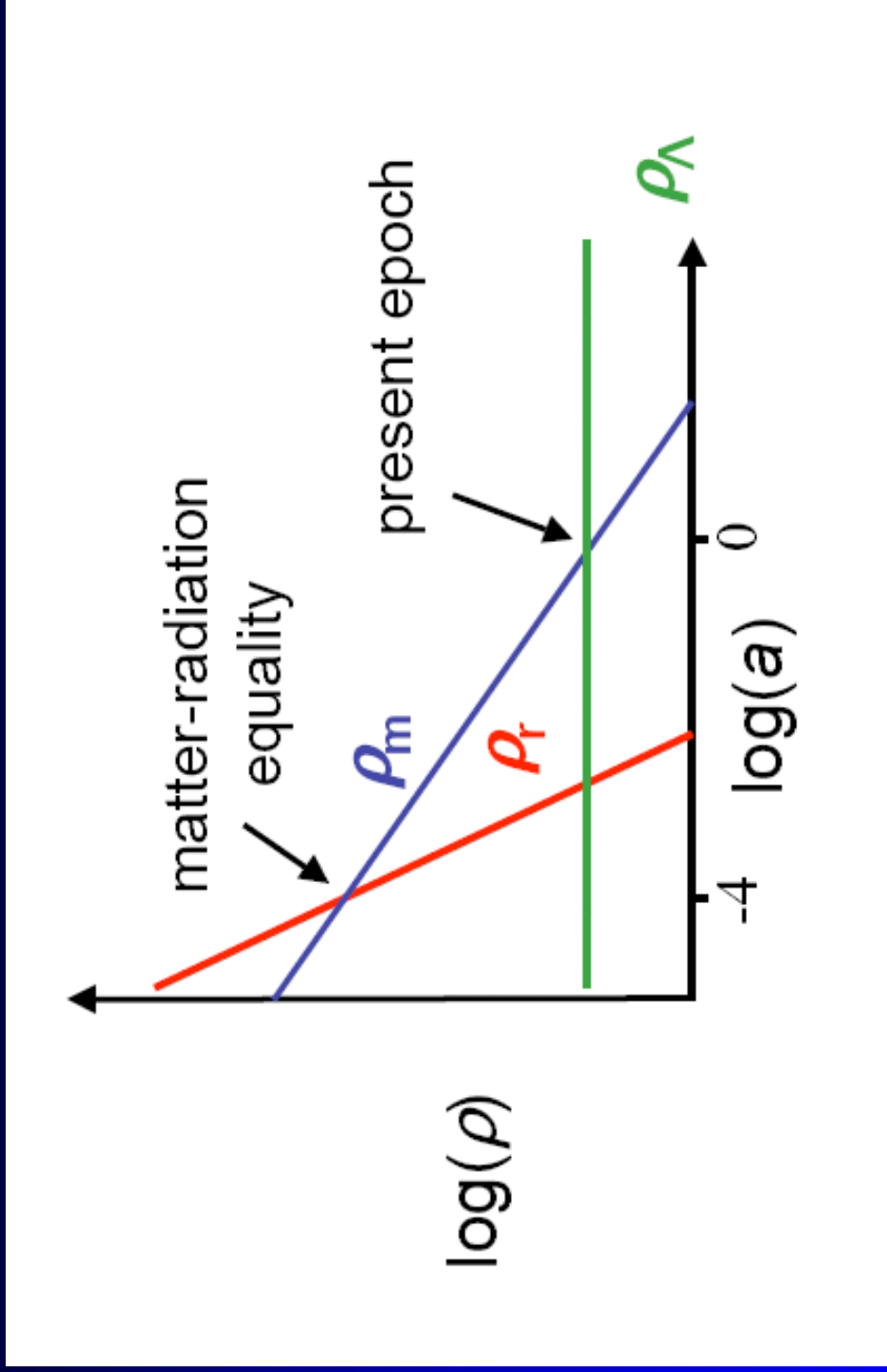
$$\rho c^2 = \frac{g}{2\pi^2} \int_m^\infty \frac{(E^2 - m^2 c^4)^{1/2}}{\exp((E - \mu)/T) \pm 1} E^2 dE$$

$$p = \frac{g}{6\pi^2} \int_m^\infty \frac{(E^2 - m^2 c^4)^{3/2}}{\exp((E - \mu)/T) \pm 1} dE.$$

relativistic  
limit

$$E^2 - m^2 c^4 \sim E^2$$

$$p = \rho c^2 / 3.$$



$$\omega(t)$$

time dependent omega



# Time dependence of scale factor $a$

➤ In the case of  $k=0$  Friedmann equation becomes

$$\dot{a}^2 = \frac{8\pi G\rho_0}{3} a^{-(1+3\omega)}$$

$$a(t) = \left(\frac{t}{t_0}\right)^{2/3+3\omega}$$

$$\omega \neq -1$$

$$a(t) = \left(\frac{t}{t_0}\right)^{1/2}$$

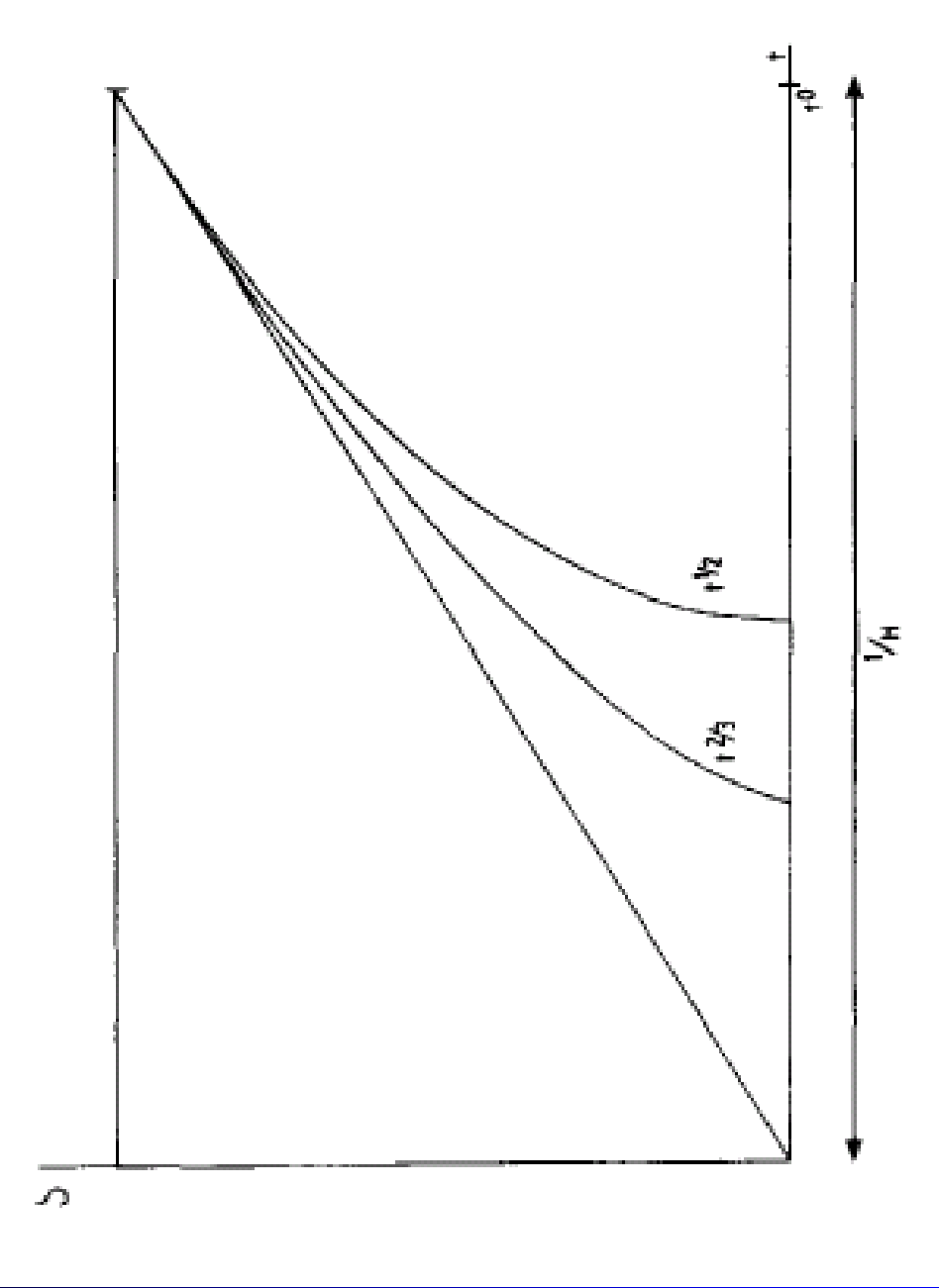
radiation  
dominated

$$a(t) = \left(\frac{t}{t_0}\right)^{2/3}$$

matter  
dominated

cosmological constant

$$a(t) \propto \exp Ht$$



In general can write

$$\frac{H^2}{H_0^2} = \frac{\Omega_{r,0}}{a^4} + \frac{\Omega_{m,0}}{a^3} + \Omega_{\Lambda,0} + \frac{1 - \Omega_0}{a^2}$$

and solve numerically

# Distances...

↗ proper distance to galaxy from which light is detected

$$d_p = c \int_{t_e}^{t_o} \frac{dt}{a(t)}$$

$$ds^2 = c^2 dt^2 - a^2(t) \left[ dr^2 + \left. \begin{array}{l} R^2 \sin^2(r/R) \\ R^2 \sinh^2(r/R) \end{array} \right\} (d\theta^2 + \sin^2 \theta d\phi^2) \right]$$

# Distances...

↗ proper distance to galaxy from which light is detected

$$c \int_{t_e}^{t_0} \frac{dt}{a(t)} = \int_0^{d_p} \frac{dr}{1 - kr^2}$$

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu = c^2 dt^2 - a^2(t) \left[ \frac{dr^2}{1 - kr^2} + r^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right]$$

# Proper distance for matter domination

$$a_0 d_p = \frac{c}{H_0 q_0^2 (1+z)} \left[ zq_0 + (q_0 - 1) \left( \sqrt{2q_0z + 1} - 1 \right) \right]$$

$$q_0 = \frac{-\ddot{a}_0}{a_0 H_0^2}$$

# Luminosity Distance

➤ In a static Euclidean space the flux detected at a distance  $d$  and the luminosity are related by

$$f = \frac{L}{4\pi d^2}$$

# Luminosity Distance

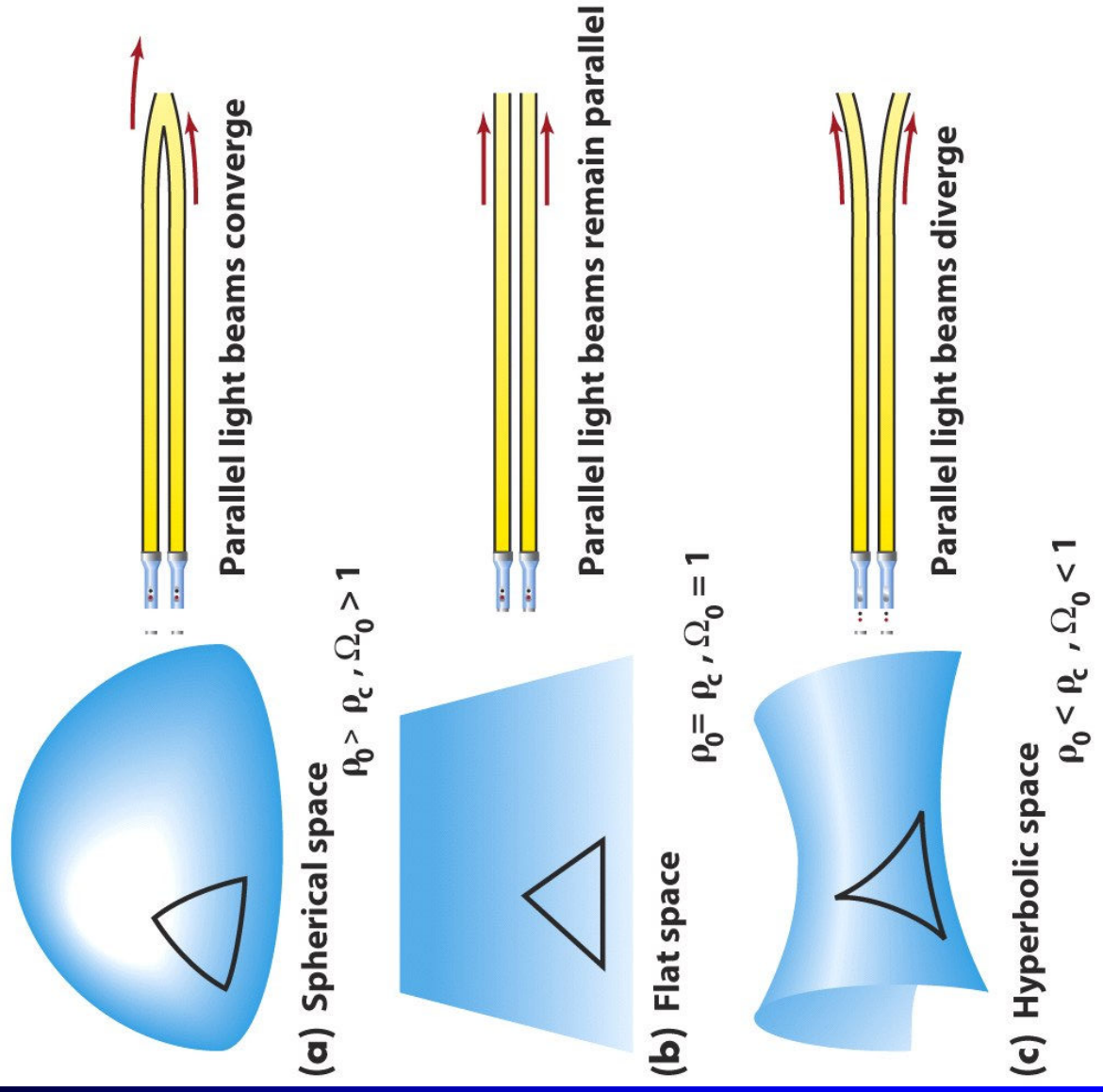
- We see light that was emitted from  $(r, \theta, \varphi)$
- The light is now spread over a sphere of proper radius

$$d_p(t_0) = r$$

and proper surface area

$$A_p(t_0) = 4\pi S_k(r)^2 = 4\pi \begin{cases} R^2 \sin^2(r/R) \\ r^2 \\ R^2 \sinh^2(r/R) \end{cases}$$





# Luminosity Distance

➤ In an expanding universe with curvature  $k$

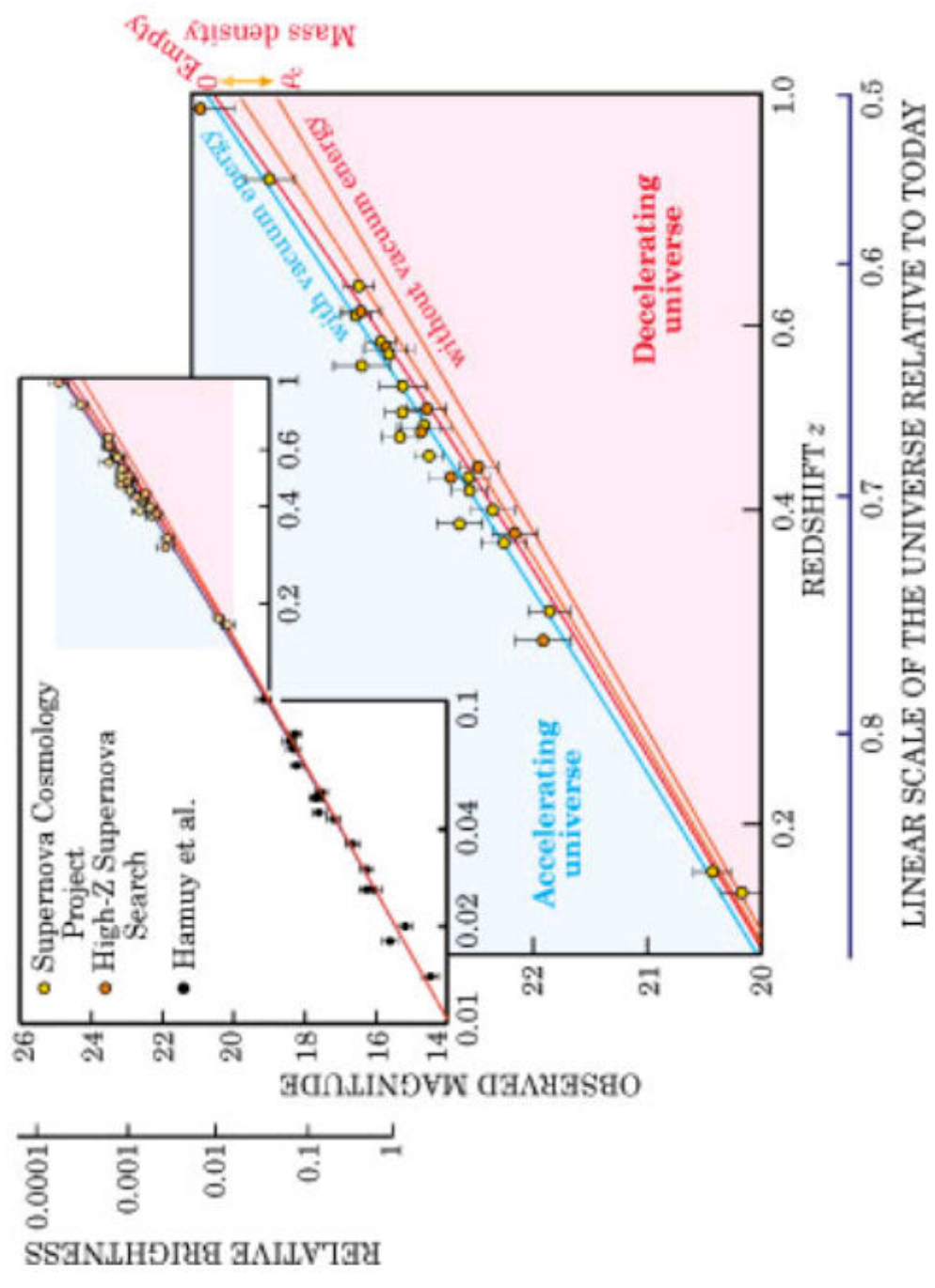
$$f = \frac{L}{4\pi S_k(r)^2 (1+z)^2}$$

➤ What is the origin of the two factors of  $(1+z)$  ?

$k=0$

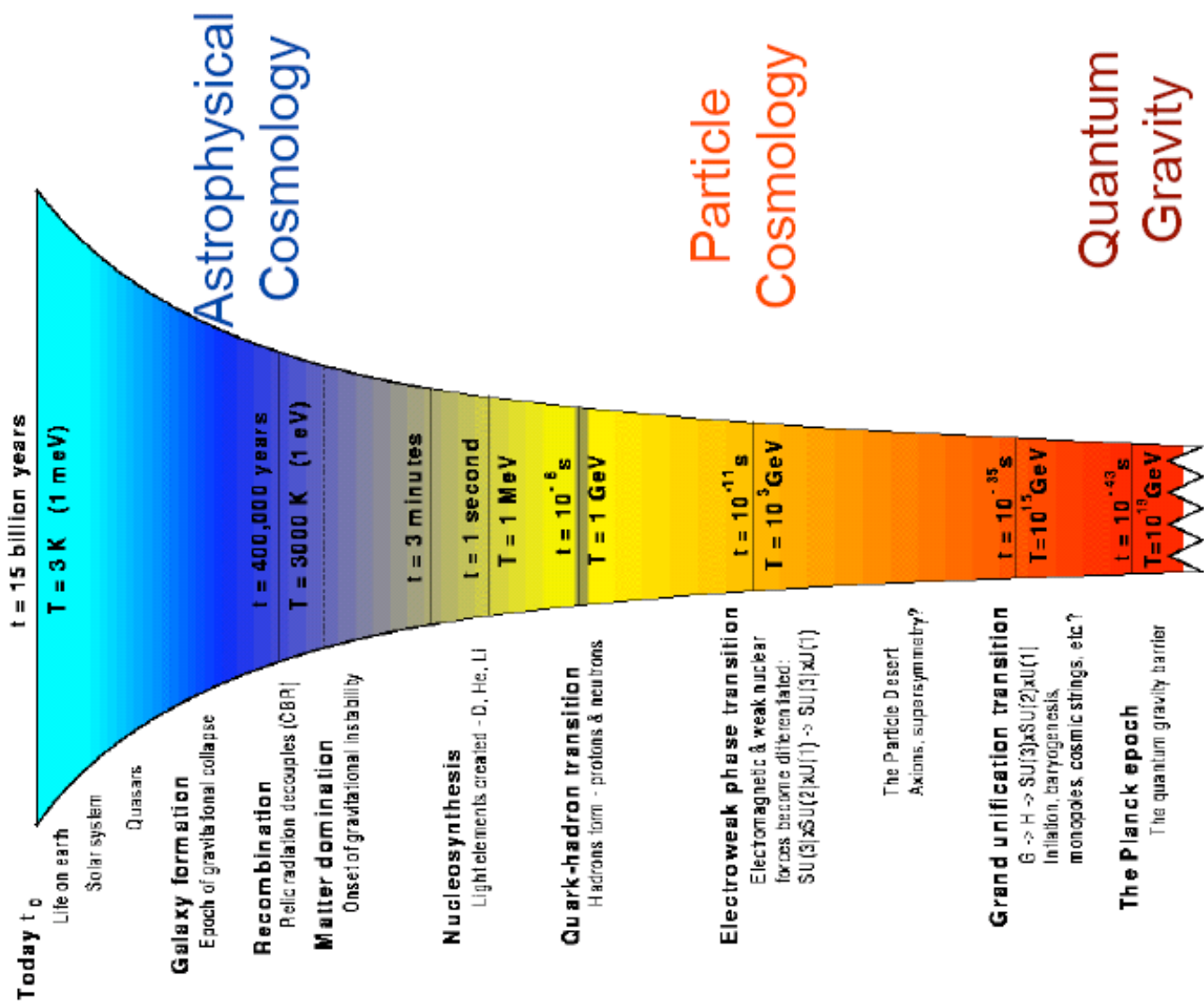
$$d_L = d_p(t_0)(1+z)$$

# Standard Candles – Type 1a Supernovae



$$M - m = -5 \log_{10} \left( \frac{d_L}{10pc} \right)$$





# Outstanding issues in the Standard Model of Cosmology...?

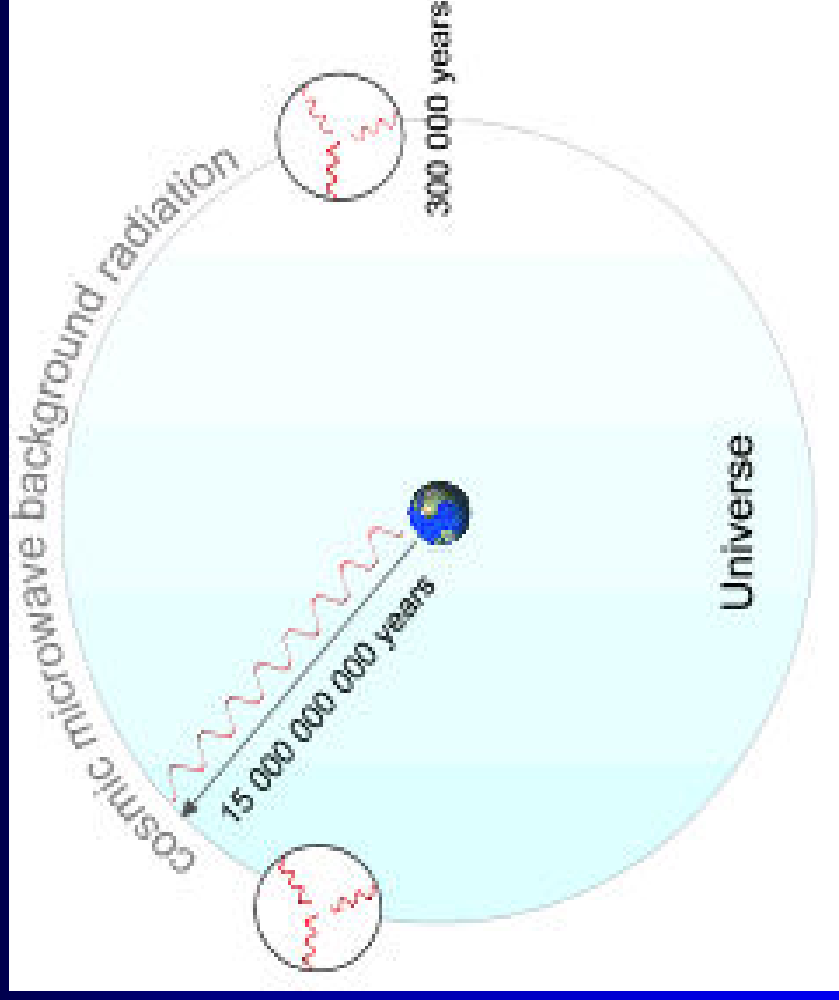
- Problems which Inflation solves
  - Horizon problem
  - Flatness problem
  - Monopole problem
  - Origin of structure seeds
- Baryon Asymmetry
- Nature of dark matter and dark energy

# Problems with the Big Bang: *Horizon Problem*

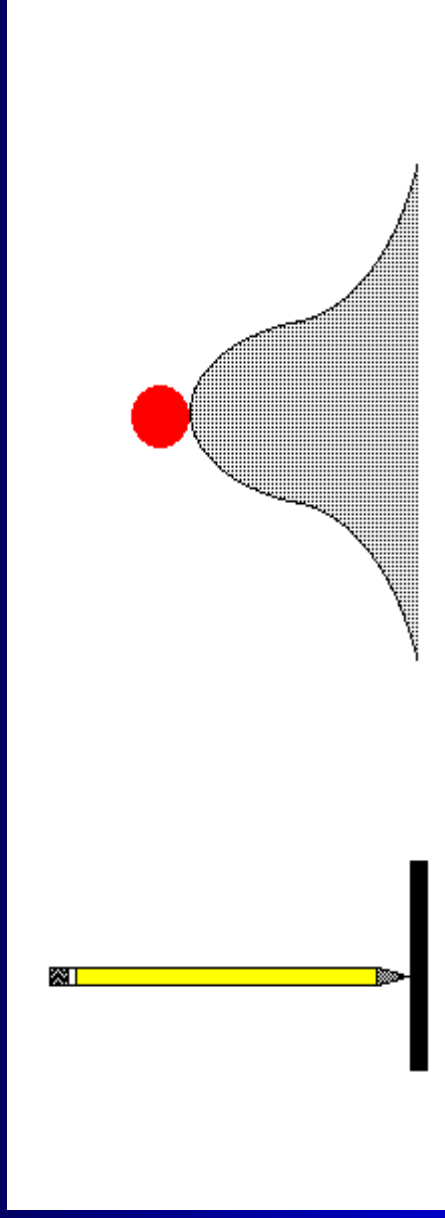
➤ The particle horizon at the time of decoupling was around 0.4 Mpc

➤ This is around  $2^\circ$  on the sky

- How can 50, 000 regions which weren't causally connected at the time of last scattering all have the same temperature (to one part in 100 000)?



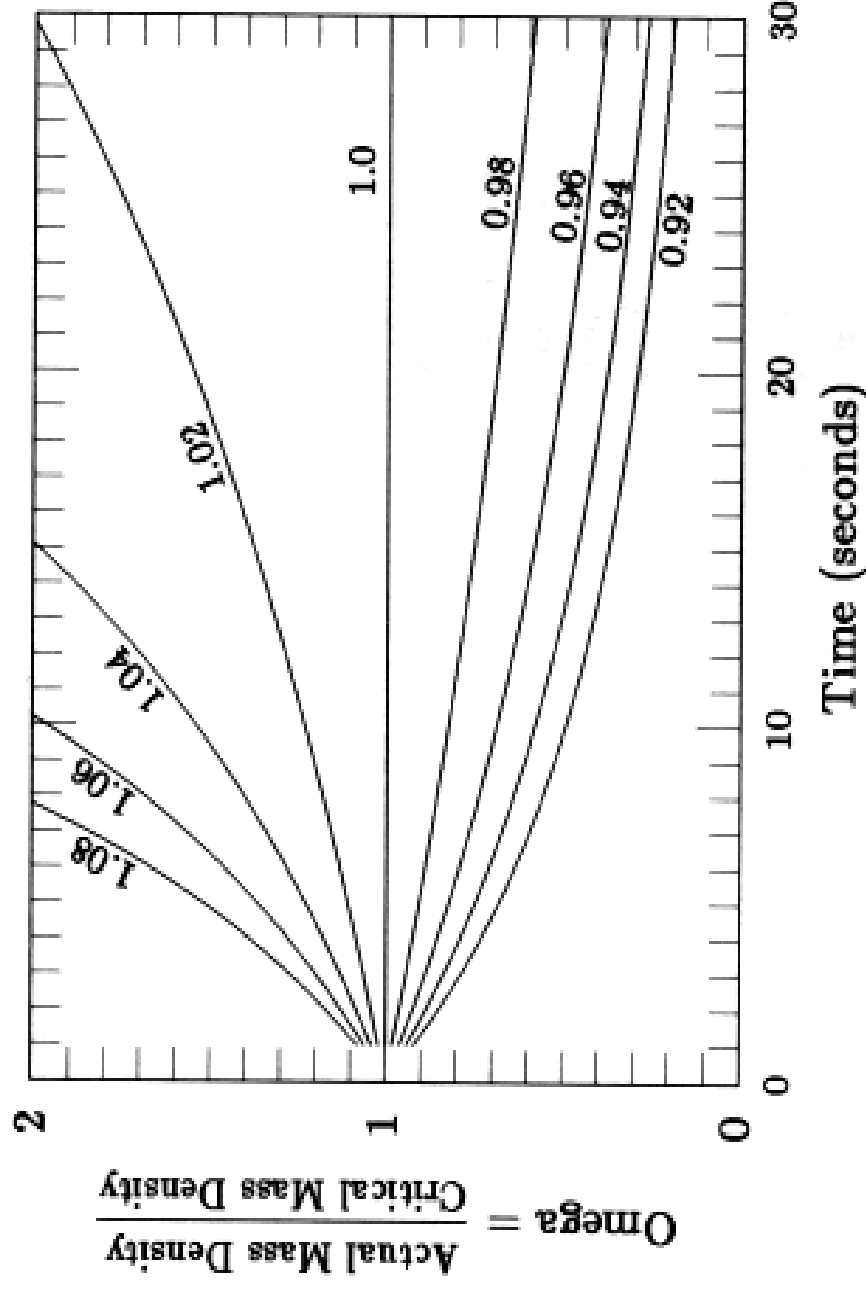
# Problems with the Big Bang: *Flatness problem*



- $\Omega = 1$  is an unstable point for the evolution
- For  $\Omega$  to be between 0.1 and 2 now it must have been extremely close to 1 in the past

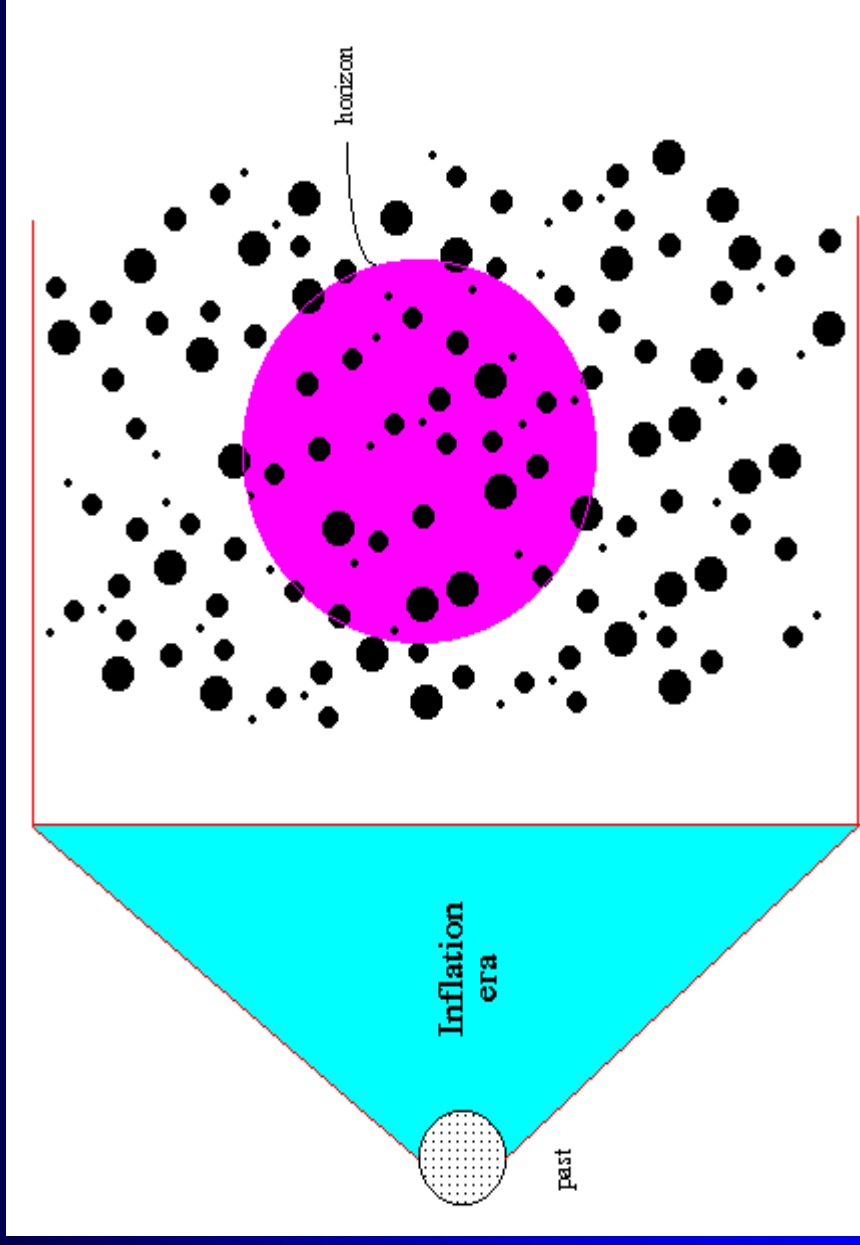


## Problems with the Big Bang: Flatness problem



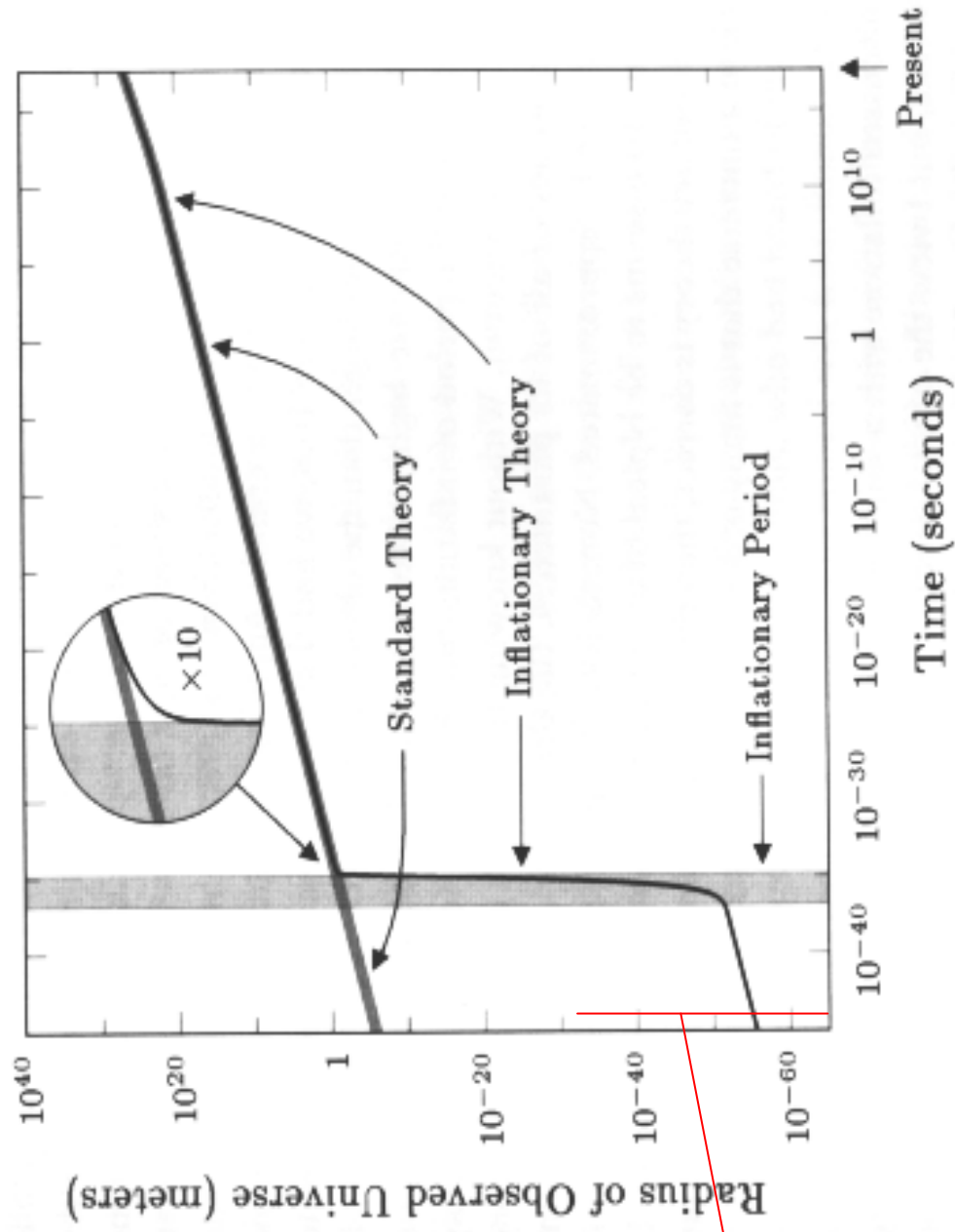
- For  $\Omega$  to be between 0.1 and 2 now it must have been extremely close to 1 in the past

# Inflation solution for: *Horizon Problem*



- If we incorporate an exponential expansion period into our calculation of the horizon at the time of last scattering, then the horizon becomes larger than the region we detect the cmb from

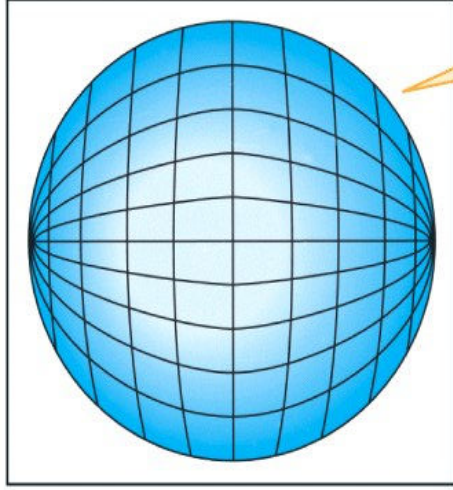
# Inflation solution for: *Horizon Problem*



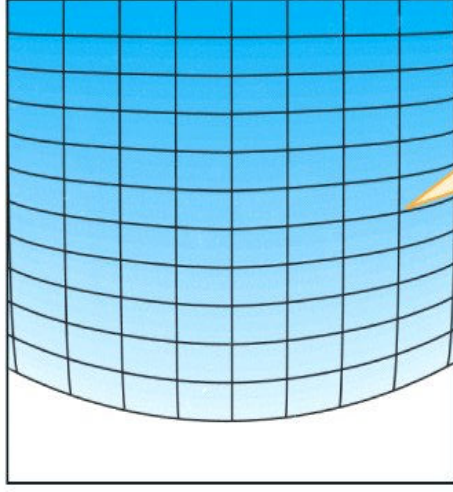
light has  
travelled this  
distance

# Inflation solution for: *Flatness problem*

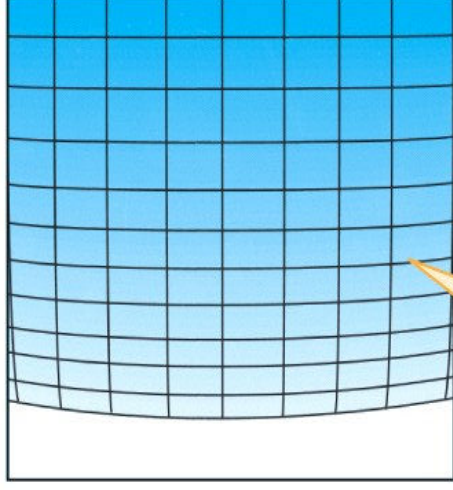
Original



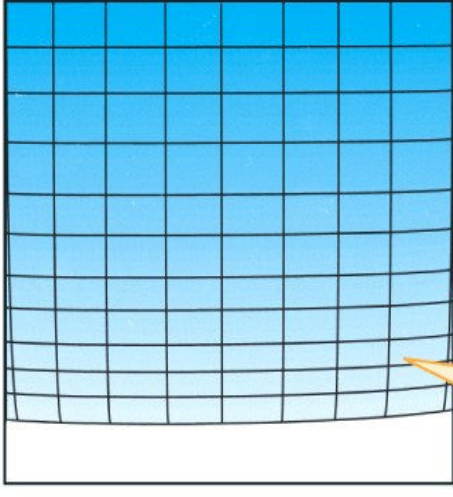
Inflated by a factor of 3 ...



by a factor of 9 ...



... and by a factor of 27.



As the sphere is inflated, its curvature eventually becomes undetectable and its surface appears flat.

- A period of inflation drives  $\Omega$  very, very close to 1 even if it was far from 1 before inflation

# Scalar field

➤ The pressure and density of a scalar field  $\phi$  are

$$\rho = \frac{1}{2} \frac{1}{\hbar c^5} \dot{\phi}^2 + \frac{V(\phi)}{c^2}$$

and

$$P = \frac{1}{2} \frac{1}{\hbar c^3} \dot{\phi}^2 - V(\phi)$$

➤ If

$$\dot{\phi} \sim 0 \quad \text{then} \quad \frac{P}{c^2} = -\rho$$

- So inflation would occur if a scalar field (generically called inflaton) had a non-zero potential energy and only slowly changing
- Inflation is not happening now so the scalar field must be at its minimum

# The inflaton rolls slowly down...

- The inflaton wants to roll down to its true vacuum, i.e. the energy minimum
- While you roll down you release energy by transforming potential energy into kinetic energy



# Baryon Asymmetry

↗ Freeze-out calculation for nucleons:

$$\text{'freeze-out' at } T \sim m_N/45, \text{ with: } \frac{n_N}{n_\gamma} = \frac{n_{\bar{N}}}{n_\gamma} \sim 10^{-19}$$

However the observed ratio is  $10^9$  times *bigger* for baryons, and *no* antibaryons are present, so there must have been an **initial asymmetry** of:

$$\frac{n_B - n_{\bar{B}}}{n_B + n_{\bar{B}}} \sim 10^{-9}$$

i.e. for every  $10^9$  baryon-antibaryon pairs there was *1 extra baryon*

# Conditions for creating baryon asymmetry

1. Baryon number violation
2. *C* and *CP* violation
3. Departure from thermal equilibrium

↗ Why not at GUT transition?



COSMOLOGY MARCHES ON

