

Curvature Perturbations from Supersymmetric Flat Directions

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A. Riotto and F.R., PLB, [arXiv:0806.3382]

Fluctuations

Dynamics

Decay

Remarks and
Conclusions II

Outline

Curvature
Perturbations
from
Supersymmetric
Flat Directions

Fluctuations During Inflation

Dynamics after inflation

Decay and Production of Perturbations

Fluctuations

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Conclusions II

Fluctuations

During Inflation:

Recall FD potential for Flat Direction *QQQL* or *uude*

$$V(\phi) = -c_I H_I^2 |\phi|^2 + \left(\lambda \frac{a_I H_I}{4M_p} \phi^4 + \text{h.c.} \right) + |\lambda|^2 \frac{|\phi|^6}{M^2},$$

- ▶ FD VEV amplitude becomes large

$$\phi_I = |\phi_I| e^{i\theta_I}$$

- ▶ Mass of θ

$$m_\theta^2 \approx 4a_I H_I^2 \ll H_I^2$$

⇒ Flat Direction phase θ_I fluctuates during inflation

$$|\delta\theta(k)|^2 \approx \frac{H_I^2}{2k^3 |\phi_I|^2}$$

⇒ Different patches have different initial conditions

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Fluctuations

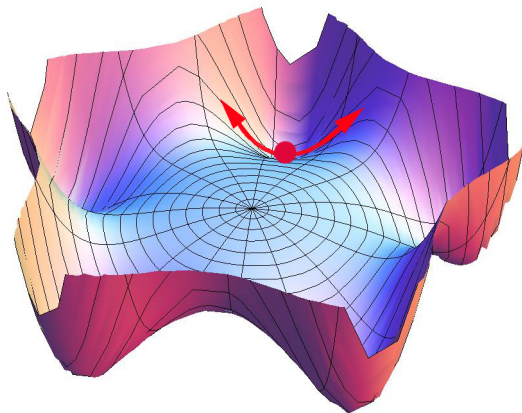
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FD Potential During Inflation

Curvature
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Flat Directions



Fluctuations

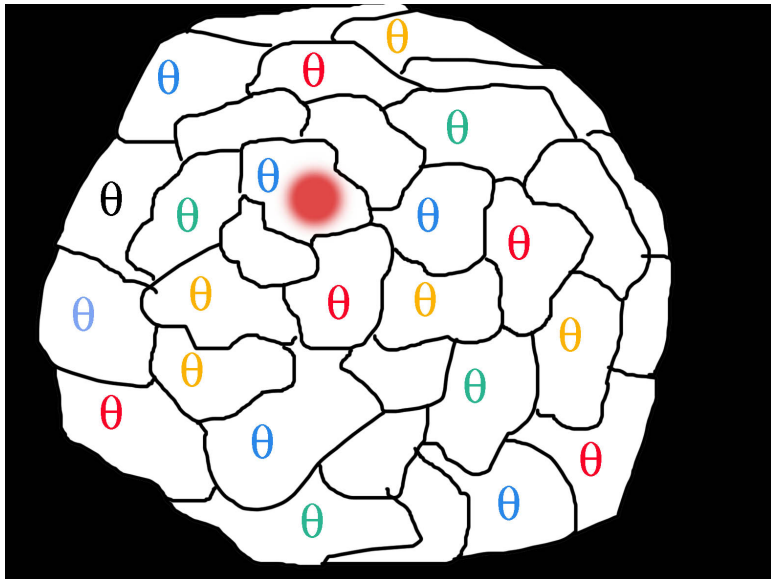
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Fluctuations in the FD Phase

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Fluctuations

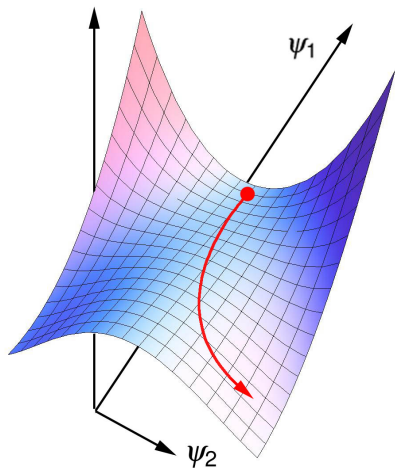
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AFTER Inflation

Curvature
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from
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Fluctuations

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After (hybrid) Inflation:

- ▶ Energy transferred from ψ_1 to ψ_2 , inflation ends
- ▶ ψ_1 and ψ_2 have different couplings to FD
- ▶ FD potential $V''(0)$ changes sign

$$V = +\frac{1}{2}H_I^2|\phi|^2 + \left(\frac{a_{\text{osc}}H_I}{M_p}\phi^4 + h.c. \right) + \lambda^2\frac{|\phi|^6}{M^2},$$

- ⇒ FD starts **oscillations** around $\phi = 0$
- ⇒ **Frequency** of oscillations $\sim H_I$

- ▶ For small $a_{\text{osc}} \neq 0$ **potential is phase-dependent**
 - ⇒ Different initial conditions θ_i have **different dynamics**

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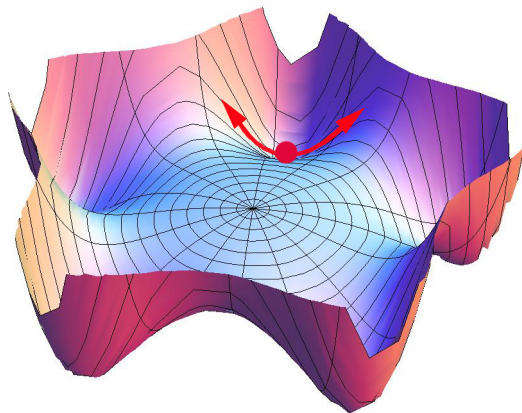
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FD Potetnial During Inflation



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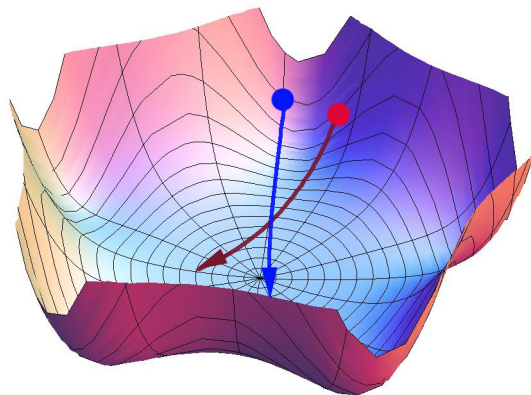
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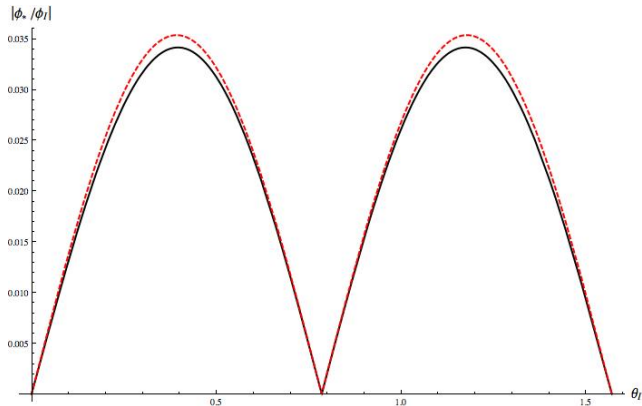
Remarks and
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Preheating

Different Trajectories

⇒ FD passes at **different distances** from origin:

$$|\phi_*| \approx \frac{2\pi}{3} a_{\text{osc}} |\phi_I| \sin(4\theta_I)$$

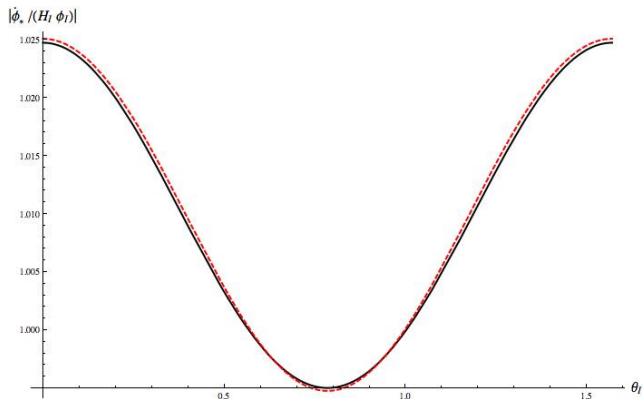


Preheating

Different Trajectories

⇒ Different speed at origin:

$$|\dot{\phi}_*| \approx H_I |\phi_I| \sqrt{1 + 4(a_{\text{osc}}/4) \cos(4\theta_I)}.$$



Preheating

...Different Distance and speed at origin:

⇒ Different decay (preheat) efficiency into light fields

$$n_{\text{light}} = \frac{(h|\dot{\phi}_*|)^{3/2}}{8\pi^3} \exp\left[-\frac{\pi h|\phi_*|^2}{|\dot{\phi}_*|}\right] \quad \text{with } \phi_* = \phi_*(\theta_I)$$

⇒ Curvature perturbation

$$\zeta = \frac{\dot{\rho}_{\text{light}}}{\dot{\rho}_{\text{tot}}} \zeta_{\text{light}} \simeq -\frac{1}{3} \frac{\rho_{\text{light}}}{\rho_{\text{tot}}} \frac{\delta n_{\text{light}}}{n_{\text{light}}},$$

⇒ Maximum spectrum for $\lambda = 0.01$

$$\mathcal{P}_\zeta^{1/2}(k) \simeq \frac{\cot(4\theta_I)}{2\pi} \left(\frac{H_I^3}{\lambda M_p^{(3)}} \right)^{1/2} \sim 2.5 \times 10^{-5},$$

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Remarks and Conclusions II

- ▶ Large and possibly dominating curvature perturbation can be produced by the non-perturbative decay of Flat Directions
- ▶ FD $QQQL$ and $uude$ don't produce baryon isocurvature perturbation
- ▶ Large Non-Gaussianities are expected