

Transient breakdown of slow-roll, homogeneity and isotropy during inflation

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UniverseNet, Oxford 22nd - 25th September 2008

Outline

- Features and motivations
- Our model
- Mode equations and initial conditions
- Results and discussion

Features in the inflaton potential

- Data from WMAP consistent with a smooth inflaton potential, with $P_{\mathcal{R}} \propto k^{n-1}$ ($n \simeq 1$)
- However 5 year WMAP data also consistent with features in the potential (*Joy, Shafieloo, Sahni and Starobinsky arXiv:0807.3334*)
- Many examples of features have been studied such as kinks, steps and bumps (*eg Adams, Cresswell and Easter arXiv:0102236 and Covi et al arXiv:0606452*)
- We wish to consider the generic effects a space dependent inflaton potential has on the primordial power spectrum

Motivations

A *phenomenological model* for the transient breakdown of slow-roll, homogeneity and isotropy during inflation

But also motivated by

- **Inhomogeneous cosmological phase transition (tachyonic preheating)**

A second field coupled to the inflaton undergoes a phase transition; A ‘mini-waterfall’ transition

- **Modulated fluctuations**

A third field imprints super-horizon inhomogeneity during a phase transition

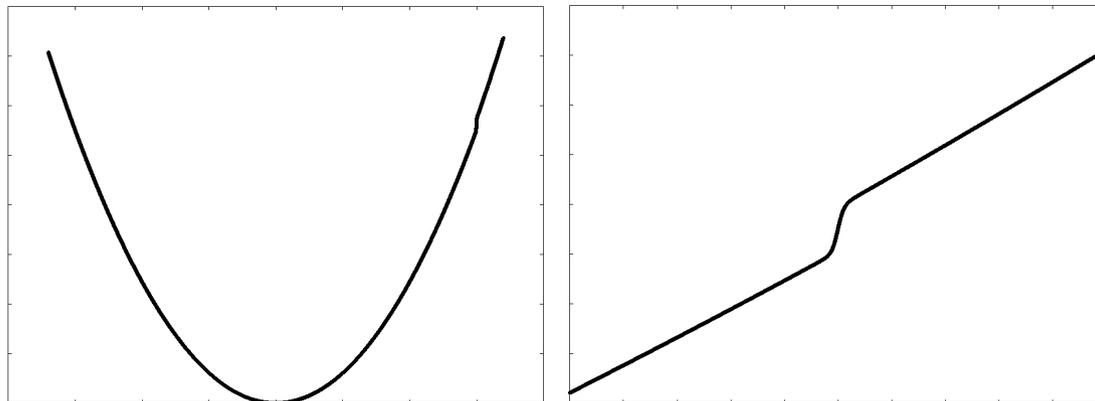
Inhomogeneity: our model

Effective potential:

$$V(\phi) = \frac{m^2 \phi^2}{2} \left(1 + c \tanh(w) \left[1 + e^{-w^2} \sin(k_L x) \sin(k_L y) \sin(k_L z) \right] \right)$$

where k_L is the inhomogeneity wavenumber and $w = \left(\frac{\phi - b}{d} \right)$

- choose b so step at $N = 55$
- solve for planar case $\sin(k_L x)$



Mode equations (1)

Start with the action $S = \frac{1}{2} \int \left(\partial_\mu u \partial^\mu u + \frac{z''}{z} u^2 \right) d^4x$,

with $u = -z\mathcal{R}$ and $z = a\dot{\phi}/H$ (\mathcal{R} is the curvature perturbation)

No space dependence: use classical field equation and quantize in term of $a_{\mathbf{k}}$ and $a_{\mathbf{k}}^\dagger$:

$$u(\mathbf{x}, \tau) = \int \frac{d^3k}{(2\pi)^{3/2}} \left[u_{\mathbf{k}}(\tau) e^{i\mathbf{k}\cdot\mathbf{x}} a_{\mathbf{k}} + u_{\mathbf{k}}^*(\tau) e^{-i\mathbf{k}\cdot\mathbf{x}} a_{\mathbf{k}}^\dagger \right]$$

giving

$$u_{\mathbf{k}}'' + \mathbf{k}^2 u_{\mathbf{k}} = \frac{z''}{z} u_{\mathbf{k}}$$

Mode equations (2)

For the **space dependent** potential, the mode equations become

$$u_{\mathbf{k}}'' + \mathbf{k}^2 u_{\mathbf{k}} - \frac{z''}{z} u_{\mathbf{k}} + \frac{a^2}{2i} \frac{d^2 F}{d\phi^2} (u_{\mathbf{k}-\mathbf{k}_L} - u_{\mathbf{k}+\mathbf{k}_L}) = 0$$

where $F(\phi) = \frac{cm^2}{2} \tanh\left(\frac{\phi-b}{d}\right) e^{-\frac{\phi-b}{d}} \phi^2$

- Equation only valid for $k \neq nk_L$ (as also classical contribution)

Initial conditions

Initial conditions set at $k^2 \gg a^2 H^2$ are

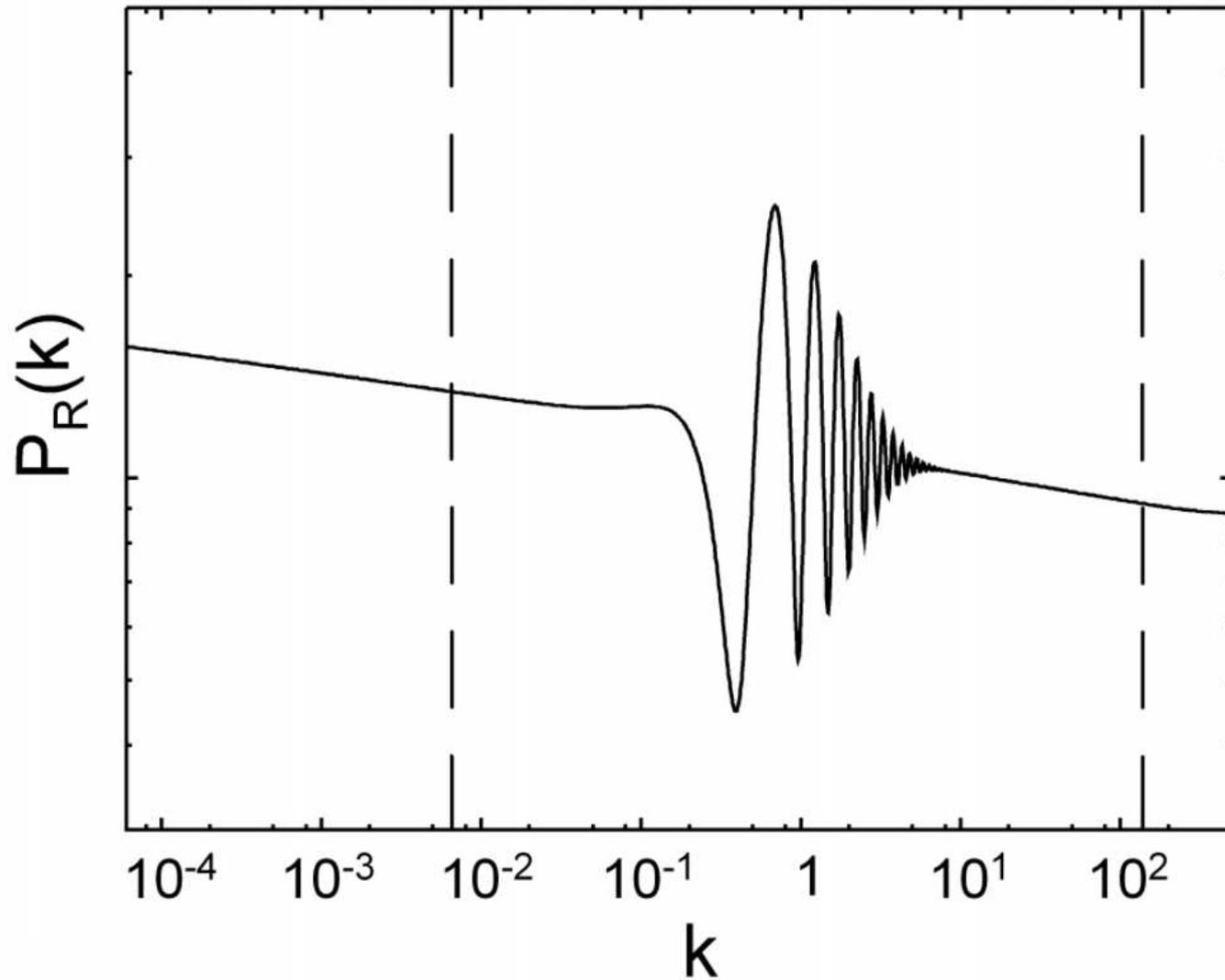
$$u_{\mathbf{k}} = \frac{e^{-i(|\mathbf{k}|\tau + \alpha_{\mathbf{k}})}}{\sqrt{2|\mathbf{k}|}}$$

where τ is conformal time and $\alpha_{\mathbf{k}}$ a random phase.
 $\alpha_{\mathbf{k}}$ only has an effect when there is mode coupling.

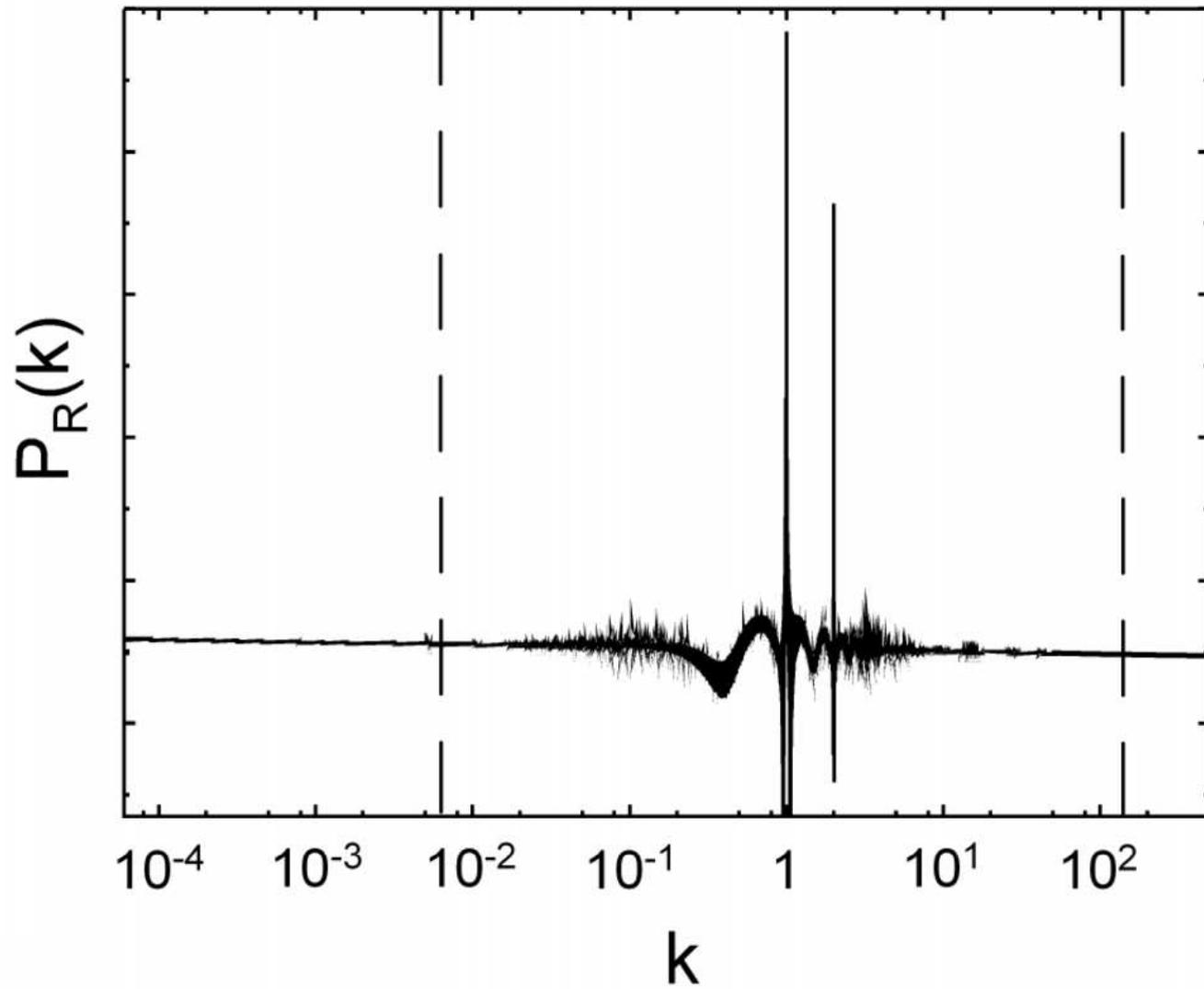
We solve for $u_{\mathbf{k}}$ and calculate the curvature perturbation power spectrum:

$$\mathcal{P}_{\mathcal{R}}^{1/2}(\mathbf{k}) = \sqrt{\frac{|\mathbf{k}|^3}{2\pi^2} \left| \frac{u_{\mathbf{k}}}{z} \right|}$$

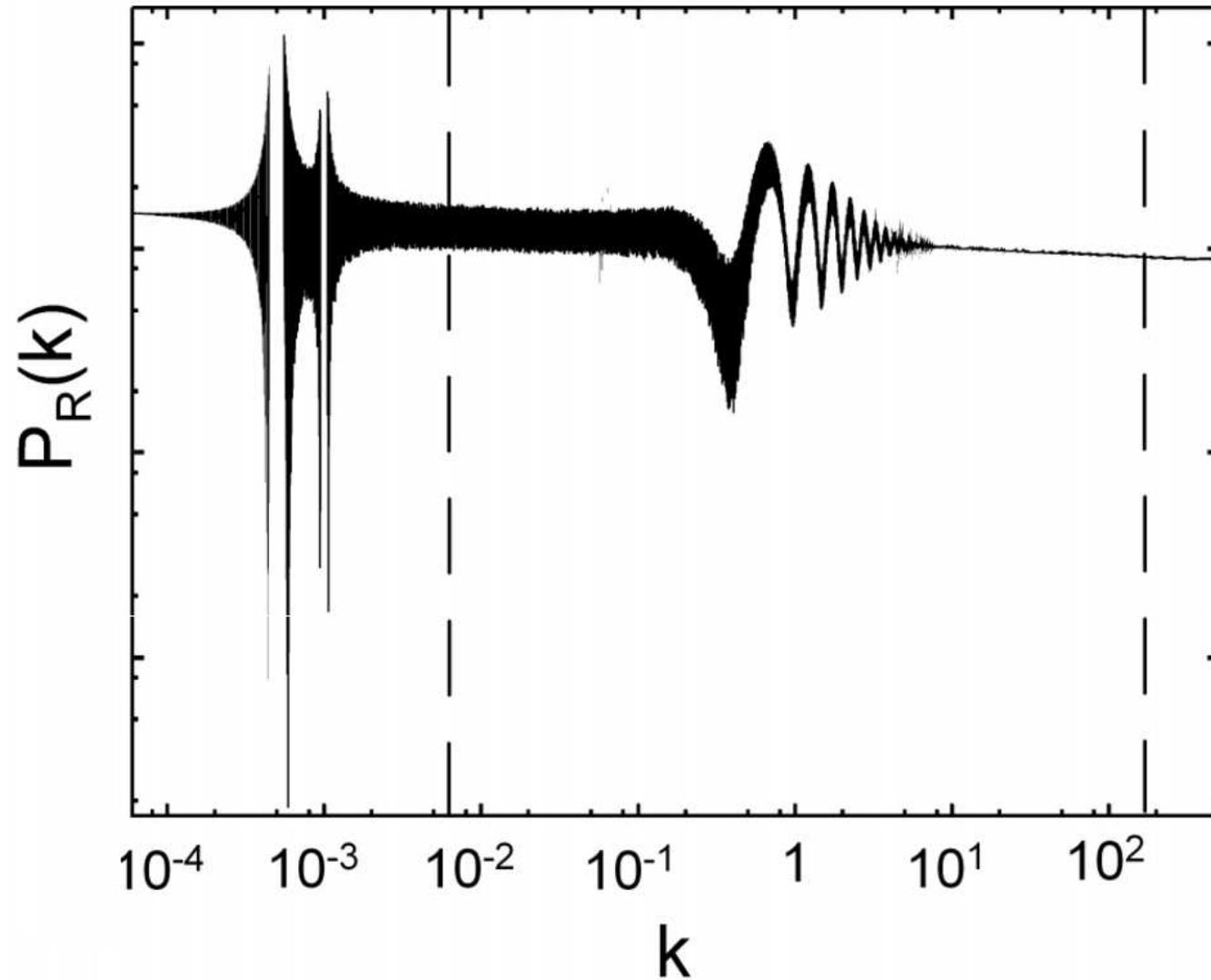
Power spectrum (homogeneous)



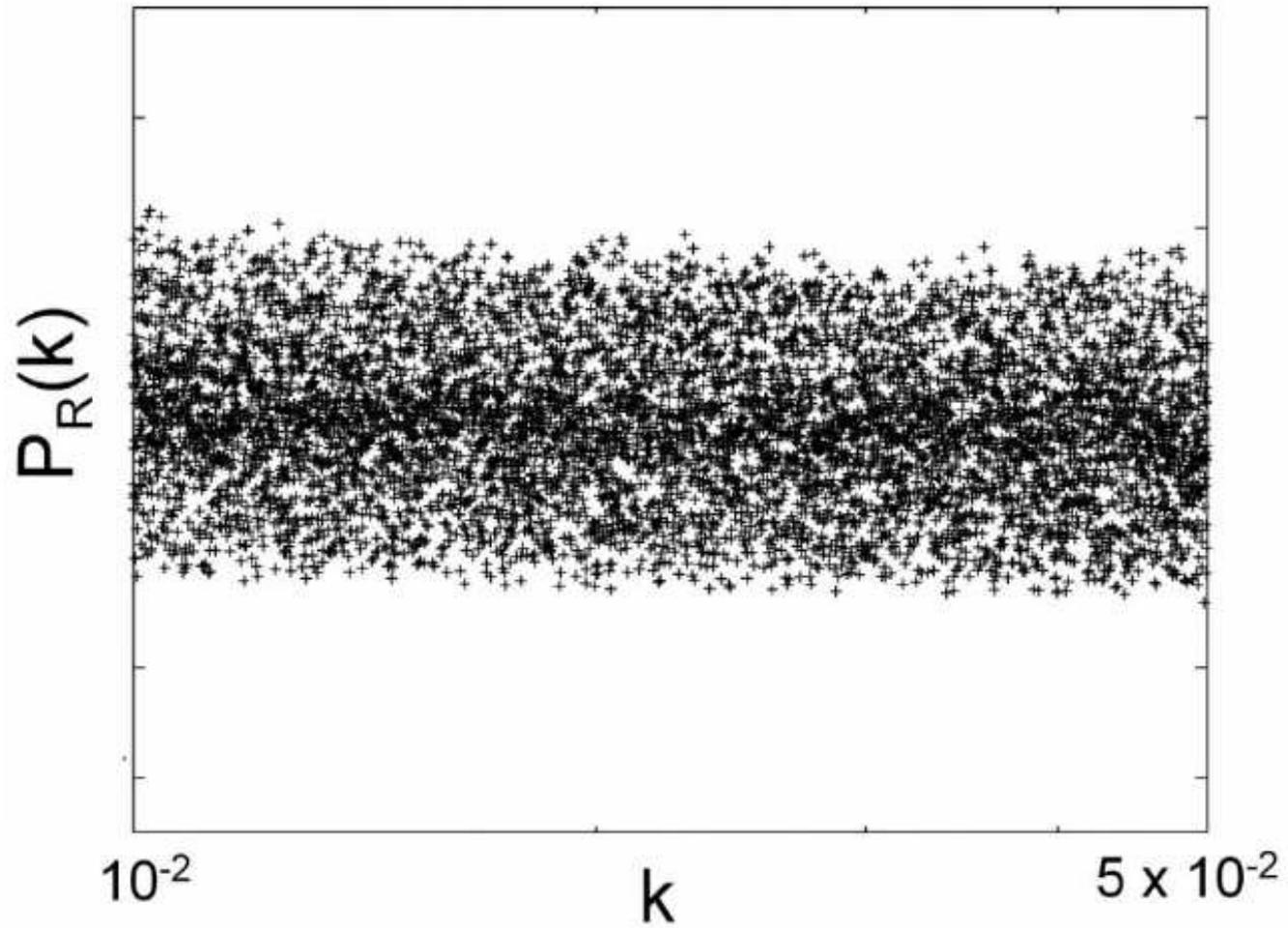
Power spectrum ($k_L = k_{STEP}$)



Power spectrum ($k_L < k_{STEP}$)



Effect of phases



Questions

- Could the resonances be observed?
- Does current data enable us to put a bound on inhomogeneity during inflation?
- Are there chances of producing **primordial black holes** at the resonances?
- Does the mode mixing generate appreciable **non-gaussian curvature perturbations**?
- Could this provide a mechanism for large scale **statistical anisotropy**?

Summary

- Transient space dependence of the inflaton potential results in **coupling** between inflaton modes
- This produces interesting results including **resonances** and random broadening
- Resonances occur when modes are coupled to amplified super-horizon modes

arXiv:coming soon!